

# 1. Brownian motion

$d$  dimensions, square lattices

basis vectors  $\hat{m}_\mu$   $\mu=1, \dots, d$

Pt on lattice:  $\vec{r} = a \sum_{\mu=1}^d m^\mu \hat{m}_\mu$   
lattice cst  $\uparrow$   $m^\mu \in \mathbb{Z}$

Random walker obeys 2 rules:

• 1: at each time step  $\Delta t$ , RW takes one step

• 2: dir<sup>n</sup> of step is uniformly distributed over all possibilities.

Need:  $P_{\vec{r}_1, t_1 | \vec{r}_0, t_0}$

conditional probab of going from  $\vec{r}_0, t_0$  to  $\vec{r}_1$  in time  $t_1 - t_0$

Sum rule:  $\sum_{\vec{r}_1} P_{\vec{r}_1, t_1 | \vec{r}_0, t_0} = 1 \quad \forall t_1$

Simple initial state:  $P_{\vec{r}_1, t_0 | \vec{r}_0, t_0} = \delta_{\vec{r}_1, \vec{r}_0}$

Discrete time evolution:

$$P_{\hat{n}_1, t_1 + \delta t | \hat{n}_0, t_0} = \frac{1}{2d} \sum_{\sigma=\pm} \sum_{m=1}^d P_{\hat{n} + a\sigma\hat{m}_m, t_1 | \hat{n}_0, t_0}$$

DTEE1

Use discrete Laplacian:

$$\nabla_a^2 f_{\hat{n}} \equiv \frac{1}{a^2} \sum_{m=1}^d \left[ f_{\hat{n} + a\hat{m}_m} + f_{\hat{n} - a\hat{m}_m} - 2f_{\hat{n}} \right]$$

$$\lim_{a \rightarrow 0} \nabla_a^2 f_{\hat{n}} = \nabla^2 f(\hat{n})$$

Time evolution:

$$P_{\hat{n}_1, t_1 + \delta t | \hat{n}_0, t_0} - P_{\hat{n}_1, t_1 | \hat{n}_0, t_0} = \frac{a^2}{2d} \nabla_a^2 P_{\hat{n}_1, t_1 | \hat{n}_0, t_0}$$

In continuum limit:

because

$$\left[ \partial_t - D \nabla^2 \right] P(\hat{n}, t) = 0$$

↑

Linear scaling  $\frac{a^2}{2d\delta t} \equiv D$

Sol<sup>n</sup>: Fourier

For  $d$  lattice:

$$f_{\vec{r}} = a^d \int_{-\pi/a}^{\pi/a} \frac{d^d k}{(2\pi)^d} e^{i\vec{k} \cdot \vec{r}} f(\vec{k})$$

$$f(\vec{k}) = \sum_{\vec{r}} e^{-i\vec{k} \cdot \vec{r}} f_{\vec{r}}$$

Initial cond<sup>n</sup>:

$$P_{\vec{r}, t_0} |_{t_0} = \delta_{\vec{r}, \vec{r}_0}$$

$$\rightarrow P_{\vec{r}, t_0} |_{t_0} = e^{-i\vec{k} \cdot \vec{r}_0}$$

TEE:

$$\sum_{\vec{r}} e^{-i\vec{k} \cdot \vec{r}} P_{\vec{r}, t_1 + \delta t} = \sum_{\vec{r}} e^{-i\vec{k} \cdot \vec{r}} \frac{1}{a^d} \sum_{\sigma=\pm} \sum_{m=1}^d P_{\vec{r} + a\sigma \hat{m}_m, t_1}$$

$$\begin{aligned} \rightarrow P_{\vec{r}, t_1 + \delta t} &= \frac{1}{a^d} \sum_{m=1}^d \sum_{\sigma=\pm} \sum_{\vec{r}'} e^{-i\vec{k} \cdot (\vec{r} - a\sigma \hat{m}_m)} P_{\vec{r}', t_1} \\ &= \left[ \frac{1}{d} \sum_{m=1}^d \cos(k_m a) \right] P_{\vec{r}, t_1} \end{aligned}$$

This is of form

$$P_{\vec{r}, t_1 + \delta t} = [\text{sum #}] P_{\vec{r}, t_1}$$

(N.B.:  $k_m = k_0 \cdot \hat{m}_m$ )

$$S_{\delta t} \therefore P_{\vec{r}, t_1} = \left[ \frac{1}{d} \sum_m \cos k_m a \right]^{\frac{t_1 - t_0}{\delta t}} P_{\vec{r}, t_0}$$

no approx<sup>n</sup>! Exact.

FT basis  $\mathbb{R}^d \simeq$  space:

$$P_{\substack{\sim_1, z_1 \\ \sim_0, z_0}} = a^d \int_{-\pi/a}^{\pi/a} \frac{d^d k}{(2\pi)^d} e^{i k \cdot (z_1 - z_0)} \left[ \frac{1}{d} \sum_{m=1}^d \cos k^m a \right]^{\frac{t_1 - t_0}{\delta t}}$$

Continuum limit:  $\frac{t_1 - t_0}{\delta t} \rightarrow$  very large

$\rightarrow$  keep only moments  $k^m$  near 0

$$\left[ \frac{1}{d} \sum_m \cos k^m a \right]^{\frac{t_1 - t_0}{\delta t}} = \left[ 1 - \frac{a^2}{2d} k^2 + \dots \right]^{\frac{t_1 - t_0}{\delta t}} \rightarrow e^{-\frac{(t_1 - t_0) a^2}{2d \delta t} k^2}$$

$$\lim_{m \rightarrow \infty} \left[ 1 + \frac{c}{m} \right]^m = e^c \quad \blacktriangleright$$

$$\left[ 1 - \frac{\frac{a^2}{2d} k^2 \left( \frac{t_1 - t_0}{\delta t} \right)}{\frac{t_1 - t_0}{\delta t}} \right]^{\frac{t_1 - t_0}{\delta t}}$$

Scaling limit:

$$\delta t \rightarrow 0, \quad a^2 \rightarrow 0$$

s.t.  $\frac{a^2}{\delta t}$  is finite, so  $\frac{a^2}{2d \delta t} \equiv D$  finite

$t_1 - t_0$  finite

Probab per unit volume:

$$P(\sim_1, z_1 | \sim_0, z_0) = \lim_{\text{scaling}} a^{-d} P_{\sim_1, z_1 | \sim_0, z_0}$$

=

$$P(\underline{r}_1, t_1 | \underline{r}_0, t_0) = \lim_{\text{scaling}} a^{-d} P_{\underline{r}_1, t_1 | \underline{r}_0, t_0}$$

$$= \int \frac{d^d \underline{k}}{(2\pi)^d} e^{-(t_1 - t_0) D \underline{k}^2 + i \underline{k} \cdot (\underline{r}_1 - \underline{r}_0)}$$

$$= \frac{1}{[4\pi D(t_1 - t_0)]^{d/2}} \exp\left[-\frac{|\underline{r}_1 - \underline{r}_0|^2}{4D(t_1 - t_0)}\right]$$

$$(\partial_t - D \nabla^2) p(\underline{r}, t) = 0$$

$$\int d^d \underline{r} p(\underline{r}, t) = 1 \quad \forall t$$



Nice property: composition

$$P(\underline{r}_2, t_2 | \underline{r}_0, t_0) =$$

$$\int d^d \underline{r}_1 P(\underline{r}_2, t_2 | \underline{r}_1, t_1) P(\underline{r}_1, t_1 | \underline{r}_0, t_0)$$

$t_2 > t_1 > t_0$

Proof: exercise.

Raw events: what is probab of  
diffusing a distance  $\propto t$ ?

Simple  $q^M$ : how much time does RW spends on a given pt  $\vec{r}_1$ ?

$$\sum_{n=0}^{\infty} P_{\vec{r}_1, t_0 + n\Delta t | \vec{r}_0, t_0} \equiv \mathcal{G}_{\vec{r}_1 - \vec{r}_0}$$

$$\sum_{n=0}^{\infty} \int_{-\pi/a}^{\pi/a} \frac{d^d k}{(2\pi)^d} e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_0)} \left[ \frac{1}{\lambda} \sum_{\vec{n}} \cos \vec{k} \cdot \vec{n} \right]^n$$

$$\sum_{n=0}^{\infty} \alpha^n = \frac{1}{1-\alpha}$$

$$\text{so } \mathcal{G}_{\vec{r}_1 - \vec{r}_0} = \int_{-\pi/a}^{\pi/a} \frac{d^d k}{(2\pi)^d} \frac{e^{i\vec{k} \cdot (\vec{r}_1 - \vec{r}_0)}}{1 - \frac{1}{\lambda} \sum_{\vec{n}} \cos \vec{k} \cdot \vec{n}}$$

We'll call this the Green's  $\mathcal{G}^M$ .

It obeys the  $q^M$

$$\mathcal{G}_{\vec{r}_1 - \vec{r}_0} = \int_{\vec{r}_1, \vec{r}_0} + \frac{1}{2d} \sum_{\vec{n}} \left[ \mathcal{G}_{\vec{r}_1 + a\vec{n}, \vec{r}_0} + \mathcal{G}_{\vec{r}_1 - a\vec{n}, \vec{r}_0} \right]$$

$$\Rightarrow -\nabla_a^2 \mathcal{G}_{\vec{r}_1 - \vec{r}_0} = \frac{2d}{a^2} \int_{\vec{r}_1, \vec{r}_0}$$

Divergences.

Focusing on region  $|\vec{r}| \sim 0$ :

$$\int d^d k \frac{e^{i\vec{k} \cdot \vec{r}}}{k^2}$$

If  $d > 2$ :

For  $d=2$ :  $\int_0^{2\pi} d\theta \int_0^{\infty} dk k \frac{e^{i\vec{k} \cdot \vec{r}}}{k^2} \sim \text{log diverg.}$

1<sup>st</sup> way to deal with div: take it away

Define a 'subtracted' Green's  $f^m$

$$M_{\vec{n}_1, \vec{n}_0}^s = M_{\vec{n}_1, \vec{n}_0} - M_0 = a^d \int \frac{d^d k}{(2\pi)^d} \frac{e^{i\vec{k} \cdot (\vec{n}_1 - \vec{n}_0)}}{1 - \frac{1}{a} \sum \cos k_\mu}$$

Ex: in 1d,  $M_{n_1, n_0}^s = -\frac{|n_1 - n_0|}{a}$

2<sup>nd</sup> way: regularization

- Rule 3: at each time interval, RW quits the game with probab  $\frac{1}{2}$

Sol<sup>m</sup>: earlier sol<sup>m</sup>, dressed with factor

$$[1 - \frac{1}{2}]^{\frac{t_1 - t_0}{\delta t}}$$

$$M_{\vec{n}_1, \vec{n}_0}^s = \sum_{n=0}^{\infty} (1 - \frac{1}{2})^n P_{\vec{n}_1, \vec{n}_0}^{t_1, t_0, \delta t}$$

$$M = \int d^d k \frac{e^{i\vec{k} \cdot \vec{n}}}{1 - \frac{(1-\frac{1}{2})}{a} \sum \cos k_\mu}$$

obey the eq<sup>m</sup>:

$$[-\nabla_a^2 + m^2] M_{\vec{n}_1, \vec{n}_0} = \frac{2d}{a^2(1-\frac{1}{2})} \delta_{\vec{n}_1, \vec{n}_0}$$

$$\frac{2d}{a^2} \frac{1}{1-\frac{1}{2}}$$

Continuous limit:  $[-\nabla^2 + m^2] M_{\vec{n}_1, \vec{n}_0} = \delta(\vec{n}_1 - \vec{n}_0)$

$$\text{sol}^m M_m(\vec{n}) = \int d^d k \frac{e^{i\vec{k} \cdot \vec{n}}}{k^2 + m^2}$$

Explicit sol<sup>ns</sup> for  $d=1,2,3$ : see notes

$$d=1: \quad \mathcal{A}_m(\nu) = \frac{e^{-m|\nu|}}{2m\nu}$$

$d=2$ : see notes

$$d=3: \quad \mathcal{A}_m(\nu) = \frac{e^{-m|\nu|}}{4\pi|\nu|}$$

towards a path integral

Reinterpretation:

$$P_{\tilde{x}_1, t_1 | \tilde{x}_0, t_0} = \frac{\# \text{ paths from } \tilde{x}_0 \text{ to } \tilde{x}_1 \text{ in time } t_1 - t_0}{\text{total } \# \text{ of paths in time } t_1 - t_0}$$

In scaling limit, read

$$P(\tilde{x}_s, t_s | \tilde{x}_i, t_i) = \frac{1}{[4\pi d(t_s - t_i)]^{d/2}} e^{-\frac{|\tilde{x}_s - \tilde{x}_i|^2}{4d(t_s - t_i)}}$$

For given  $t_s$  &  $t_i$ , 'slice' the interval into  $N$  pieces (and let  $N \rightarrow \infty$  eventually)

Use composition property:  $t_n = t_i + n\Delta t$

$$P(\tilde{x}_s, t_s | \tilde{x}_i, t_i) = \int \prod_{n=1}^{N-1} d\tilde{x}_n P(\tilde{x}_s, t_s | \tilde{x}_{N-1}, t_{N-1}) \times \frac{t_s - t_i}{N} \\ \times P(\tilde{x}_{N-1}, t_{N-1} | \tilde{x}_{N-2}, t_{N-2}) \cdots P(\tilde{x}_1, t_1 | \tilde{x}_i, t_i)$$



For a small interval  $\Delta t$ , we have (letting  $\Delta t \rightarrow 0$ )

$$P(N_{m+1}, t_{m+1} | N_m, t_m) \rightarrow \frac{1}{[4\pi D \Delta t]^{\frac{d}{2}}} \exp\left\{-\frac{|N_{m+1} - N_m|^2}{4D \Delta t}\right\}$$

$$\text{Write } \frac{N_{m+1} - N_m}{\Delta t} \rightarrow \frac{dN(t)}{dt}$$

$\rightarrow$  our proba becomes a path integral:

$$P(N_s, t_s | N_i, t_i) = \int_{\substack{N(t_i) = N_i \\ N(t_s) = N_s}} \mathcal{D}N(t) \exp\left\{-\frac{1}{4D} \int_{t_i}^{t_s} dt \left|\frac{dN(t)}{dt}\right|^2\right\}$$

Path integral measure:

$\mathcal{D}$ : see eq<sup>m</sup> 77 in notes.